

Unconstrained all-order higher-spin vertices

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Higher-spin problem

$$S_{HS} = \int \sum_s \phi_s \square_F \phi_s + \sum_{s_1, s_2, s_3} a_{s_1 s_2 s_3} D..D\phi_{s_1} D..D\phi_{s_2} D..D\phi_{s_3} + O(\phi^4)$$

Quadratic action

Fronsdal, Fang and Fronsdal:

$$\delta\phi_{m_1..m_s} = D_{(m_1} \epsilon_{m_2..m_s)} , \quad g^{mn} g^{kl} \phi_{mnkl...} = 0$$

Cubic action

Bengtsson², Brink; Berends, Burgers, van Dam; Metsaev; Fradkin, Vasiliev; many more

$$s_1 + s_2 + s_3 - 2\min(s_i) \leq N(D) \leq s_1 + s_2 + s_3$$

No go: Minkowski \rightarrow AdS \rightarrow higher-spin symmetry

Higher-spin holography

Klebanov, Polyakov; (Sezgin, Sundell; Leigh, Petkou)

$$S_{HS} \Big|_{z \rightarrow 0} = O(N)\text{-model}$$

3pt test (Giombi, Yin): OK

Maldacena, Zhiboedov:

$$\langle J_{s_1} J_{s_2} J_{s_3} \rangle = \sin^2 \phi \langle J_{s_1} J_{s_2} J_{s_3} \rangle_b + \cos^2 \phi \langle J_{s_1} J_{s_2} J_{s_3} \rangle_f + \frac{1}{2} \sin 2\phi \langle J_{s_1} J_{s_2} J_{s_3} \rangle_o$$

In the bulk (Tatarenko, Vasiliev)

Holographic reconstruction (Sleight, Taronna):

$$\langle J_{s_1} J_{s_2} J_{s_3} \rangle \rightarrow S_{HS}^{(3)}$$

Locality problem

Quartic reconstruction

Bekaert, Erdmenger, Ponomarev, Sleight:

$$\langle J_{s_1} J_{s_2} J_{s_3} J_{s_4} \rangle \rightarrow S_{HS}^{(4)} \quad (\text{nonlocal})$$

Sleight, Taronna: A type of the nonlocality is

$$\frac{1}{\square} \quad (?)$$

Gelfond: Direct calculation from Vasiliev equations

(moderate nonlocality)

Locality: state of the art

- HS theory in four dimensions

$$HS_4 = \text{hol} \oplus \text{ahol} \oplus \text{mixed}$$

Hol: local up to quite higher orders (V.D., Gelfond, Korybut, Vasiliev)

- **Hol:** generating equations; all-order locality (V.D.)
- **Mixed:** Moderate nonlocality at quartic order (Gelfond)
- Bosonic HS theory in any dimension: generating equations; all-order locality at the **unconstrained** level (V.D., Korybut)

Vasiliev's unfolded approach

Stick to the frame-like formalism instead of the metric-like one:

$$\phi_{a_1..a_s} \rightarrow \begin{cases} \omega(Y|x) = dx^\mu \omega_\mu(Y|x) & - \text{potentials} \\ C(Y|x) & - \text{curvatures} \end{cases},$$

Y^J are the auxiliary generating variables

Nonlinear HS equations of motion:

$$d\omega = -\omega * \omega + \mathcal{V}(\omega, \omega, C) + \dots + \mathcal{V}(\omega, \omega, C..C),$$

$$dC = -\omega * C + C * \pi(\omega) + \Upsilon(\omega, C, C) + \dots + \Upsilon(\omega, C..C).$$

Consistency

$$d^2 = 0 \rightarrow \mathcal{V}(\omega, \omega, C..C), \quad \Upsilon(\omega, C..C) \quad (\text{Vertices})$$

$$\omega \rightarrow \omega + f_\omega(\omega, C..C), \quad C \rightarrow C + f_C(C..C) \quad (\text{field redefinition})$$

- * – associative HS algebra (spanned by Y 's)

$$f(Y) * g(Y) = f(Y) e^{i \overleftarrow{\partial}_Y \cdot \overrightarrow{\partial}_Y} g(Y)$$

- Lower order vertices

$$\mathcal{V}(\Omega, \Omega, C) \rightarrow \text{Free equations} \quad (\Omega\text{-vacuum})$$

$$\omega * \omega, \quad \omega * C - C * \pi(\omega), \quad \mathcal{V}(\Omega, \omega, C), \quad \Upsilon(\Omega, C, C) \rightarrow \text{Cubic int.}$$

- d -dimensional off-shell bosonic HS algebra:

$$\omega^{a(s-1), b(n)} = dx^\mu \omega_\mu^{a(s-1), b(n)} = \begin{array}{|c|c|c|c|c|} \hline & \bullet & \bullet & \bullet & \\ \hline & \bullet & \bullet & & \\ \hline \end{array}, \quad 0 \leq n \leq s-1$$

$$C^{a(m), b(s)} = \begin{array}{|c|c|c|c|c|} \hline & \bullet & \bullet & \bullet & \\ \hline & \bullet & \bullet & & \\ \hline \end{array}, \quad m \geq s$$

$$Y^J = (y_\alpha, y_\alpha^a), \quad \alpha = 1, 2, \quad a = 0..d-1$$

$$(f * g)(y, \vec{y}) = \int f(y + u, \vec{y} + \vec{u}) g(y + v, \vec{y} + \vec{v}) e^{iuv + i\vec{u}\vec{v}}$$

Generating interactions (unconstrained)

Off-shell form $Y = (y, \mathbf{y})$

$$d\omega = -\omega * \omega + \mathcal{V}(\omega, \omega, C) + \dots + \mathcal{V}(\omega, \omega, C..C),$$

$$dC = -\omega * C + C * \pi(\omega) + \Upsilon(\omega, C, C) + \dots + \Upsilon(\omega, C..C).$$

$$\pi f(y, \vec{\mathbf{y}}) = f(-y, \vec{\mathbf{y}})$$

Vasiliev's trick:

$$\omega \rightarrow \textcolor{red}{W}(z; Y|x) := \omega(Y|x) + W_1(z; Y|x) + W_2(z; Y|x) + \dots,$$

$$d_x W + W * W = 0$$

star product is extended to (z, Y) – space

Generating system

V.D.'22; V.D., Korybut'23:

star product:

$$(f * g)(z; Y) = \int f(z + u', y + u; \mathbf{y}) \star g(z - v, y + v + v'; \mathbf{y}) e^{iuv + iu'v'} ,$$

Equations:

$$d_x W + W * W = 0 ,$$

$$d_z W + \{W, \Lambda\}_* + d_x \Lambda = 0 ,$$

$$d_x C + (W(z'; y, \vec{\mathbf{y}}) * C - C * W(z'; -y, \vec{\mathbf{y}})) \Big|_{z'=-y} = 0$$

$$\Lambda = dz^\alpha z_\alpha \int_0^1 d\tau \tau C(-\tau z, \vec{\mathbf{y}}) e^{i\tau z y}$$

Comparison with Vasiliev

Vasiliev equations:

$$d_x W + W * W = 0,$$

$$d_x B + W * B - B * \pi(W) = 0,$$

$$d_z W + \{W, \Lambda\}_* + d_x \Lambda = 0,$$

$$d_z \Lambda + \Lambda * \Lambda = i dz^\alpha dz_\alpha B * \kappa, \quad \kappa = e^{izy}$$

$$d_z B + \Lambda * B - B * \pi(\Lambda) = 0.$$

star product:

$$(f * g)(z; Y) = \int f(z+u, y+u; \mathbf{y}) \star g(z-v, y+v; \mathbf{y}) e^{iuv},$$

Perturbative series

Canonical embedding:

$$W = \omega(Y|x) + W^{(1)}(z, Y|x) + W^{(2)}(z, Y|x) + \dots, \quad W^{(n)}(0, Y|x) = 0, \quad \forall n \geq 1$$

Vertices:

$$\mathcal{V}(\omega, \omega, C^n) = - \left(\sum_{k+m=n} W^{(k)} * W^{(m)} \right) \Big|_{z=0}$$

$$\Upsilon(\omega, C^n) = - \left(W^{(n)}(z'; y, \vec{\mathbf{y}}) * C - C * W^{(n)}(z'; -y, \vec{\mathbf{y}}) \right) \Big|_{z'=-y}$$

Source prescription (VD, Gelfond, Korybut, Vasiliev)

$$\omega(y, \vec{\mathbf{y}}) = \exp \left(y_\alpha \frac{\partial}{\partial y_\alpha^t} \right) \omega(y_\alpha^t, \vec{\mathbf{y}}) \Big|_{y_\alpha^t=0} = e^{-iy_\alpha t^\alpha} \omega(y_\alpha^t, \vec{\mathbf{y}}) \Big|_{y_\alpha^t=0}, \quad t^\alpha := i \frac{\partial}{\partial y_\alpha^t}$$

$$C(y, \vec{\mathbf{y}}) = \exp \left(y_\alpha \frac{\partial}{\partial y_\alpha^p} \right) C(y_\alpha^p, \vec{\mathbf{y}}) \Big|_{y_\alpha^p=0} = e^{-iy_\alpha p^\alpha} C(y_\alpha^p, \vec{\mathbf{y}}) \Big|_{y_\alpha^p=0}, \quad p^\alpha := i \frac{\partial}{\partial y_\alpha^p}$$

Master field W :

$$W^{(n)}(z; Y) = \sum_{k=0}^n \mathcal{W}^{(k|n)}(z; y|t; p_1..p_n) (\overbrace{C...C}^k \omega \overbrace{C...C}^{n-k}) \Big|_{y^t, y^p=0}$$

$$(\overbrace{C...C}^k \omega \overbrace{C...C}^{n-k}) = \dots \star C\left(y^{p_k}, \vec{y}\right) \star \omega\left(y^t, \vec{y}\right) \star C\left(y^{p_{k+1}}, \vec{y}\right) \star \dots$$

Vertices

$$\Upsilon(\omega, C^n) = \sum_{k=0}^n \Phi_n^{[k]}(y|t; p_1..p_n) (\overbrace{C...C}^k \omega \overbrace{C...C}^{n-k}) \Big|_{y^t, y^p=0}$$

$$\mathcal{V}(\omega, \omega, C^n) = \sum_{0 \leq k_1 \leq k_2 \leq n} \Psi_n^{[k_1, k_2]}(y|t_1, t_2; p_1..p_n) (\overbrace{C...C}^{k_1} \omega \overbrace{C...C}^{k_2 - k_1} \omega \overbrace{C...C}^{n - k_2}) \Big|_{y^t, y^p=0}$$

W calculation

Iterative equations

$$d_z W^{(n+1)} + \{W^{(n)}, \Lambda\}_* + (d_x \Lambda)^{(n+1)} = 0, \quad W^{(0)} = \omega(Y|x)$$

Standard homotopy solution

$$W^{(n+1)} = -\Delta(\{W^{(n)}, \Lambda\}_*), \quad \Delta(f_\alpha(z) dz^\alpha) := z^\alpha \int_0^1 d\tau f_\alpha(\tau z)$$

In terms of sources

$$\begin{cases} \mathcal{W}^{(k|n)} = \Delta[\mathcal{W}^{(k|n-1)} * \Lambda' - \Lambda' * \mathcal{W}^{(k-1|n-1)}], & k \in [0, n] \\ \mathcal{W}^{(0|0)} = e^{-iyt} \end{cases}$$

$$\Lambda'(z; y|p) = \cancel{dz^\alpha z_\alpha \int_0^1 d\sigma \sigma e^{i\sigma z(y+p)}} \longrightarrow \Lambda(z; Y) = \Lambda'(z; y|p) C(y^p, \vec{y}) \Big|_{y^p=0}$$

Vertex Υ

$$\Upsilon(\omega, C^n) \rightarrow \Phi_n^{[k]}(y|t; p_1, \dots, p_n) = (-)^{k+1} (ty)^{n-1} \int_{\mathcal{D}_n^{[k]}} d\xi d\eta e^{-iyP_n(\xi) - itP_n(\eta) + i(ty) \cdot S_n^{[k]}}$$

$$P_n(\zeta) = \sum_{s=1}^n \zeta_s p_s, \quad S_n^{[k]} = - \sum_{s=1}^k \xi_s + \sum_{s=k+1}^n \xi_s + \sum_{i < j} (\xi_i \eta_j - \xi_j \eta_i)$$

$$\mathcal{D}_n^{[k]} = \begin{cases} \eta_1 + \dots + \eta_n = 1, \quad \eta_i \geq 0, \\ \xi_1 + \dots + \xi_n = 1, \quad \xi_i \geq 0, \\ \eta_i \xi_{i+1} - \eta_{i+1} \xi_i \leq 0, \quad i \in [1, k-1], \\ \eta_i \xi_{i+1} - \eta_{i+1} \xi_i \geq 0, \quad i \in [k+1, n-1]. \end{cases}$$

Vertex Υ

$$\begin{aligned} \Upsilon(\omega, C^n) = & (-i)^{n-1} \sum_{k=0}^n (-)^{k+1} \int_{\mathcal{D}_n^{[k]}} d\xi d\eta \int \frac{d^2 u d^2 v}{(2\pi)^2} e^{iuv} \times \\ & \times \left(\prod_{i=1}^k \star C(\xi_i y + \eta_i v, \vec{\mathbf{y}}) \right) \star (y^\alpha \partial_\alpha)^{n-1} \omega(S_n^{[k]} y + u, \vec{\mathbf{y}}) \star \left(\prod_{j=k+1}^n \star C(\xi_j y + \eta_j v, \vec{\mathbf{y}}) \right) \end{aligned}$$

Geometric interpretation

New variables

$$x_0 = 0, \quad x_i = \sum_{s=1}^i \xi_s,$$

$$y_0 = 0, \quad y_i = \sum_{s=1}^i \eta_s$$

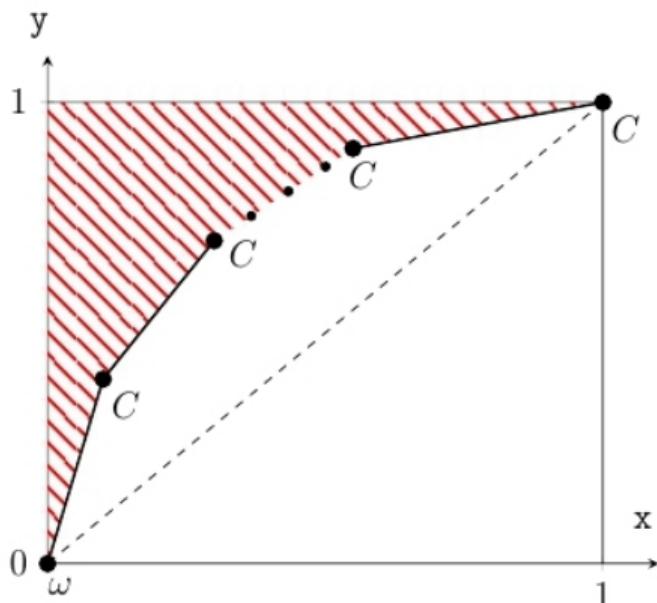
$$x_i \geq x_j, \quad y_i \geq y_j, \quad i > j$$

$$x_i \leq 1, \quad y_i \leq 1, \quad x_n = y_n = 1$$

Geometric interpretation

ω -left case:

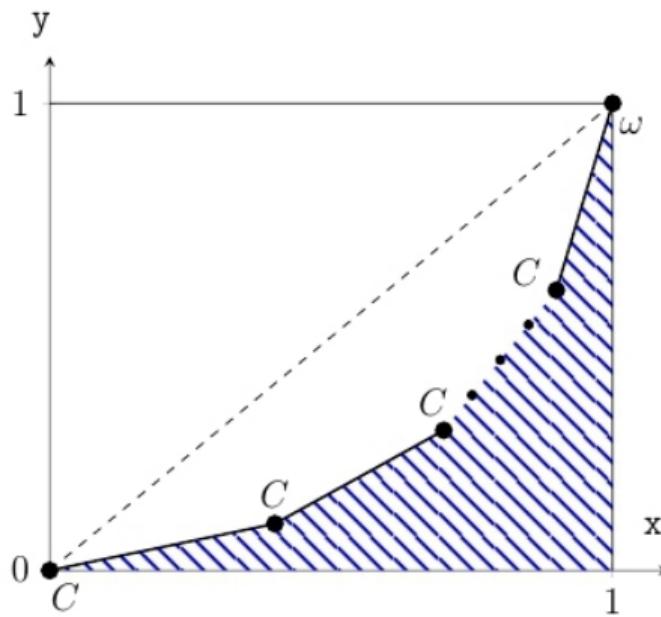
$$\mathcal{D}_n^{[0]} : \quad \eta_i \xi_{i+1} - \eta_{i+1} \xi_i \geq 0, \quad i = 1, \dots, n-1; \quad \frac{1}{2} S_n^{[0]} = \text{Area}$$



Geometric interpretation

ω -right case:

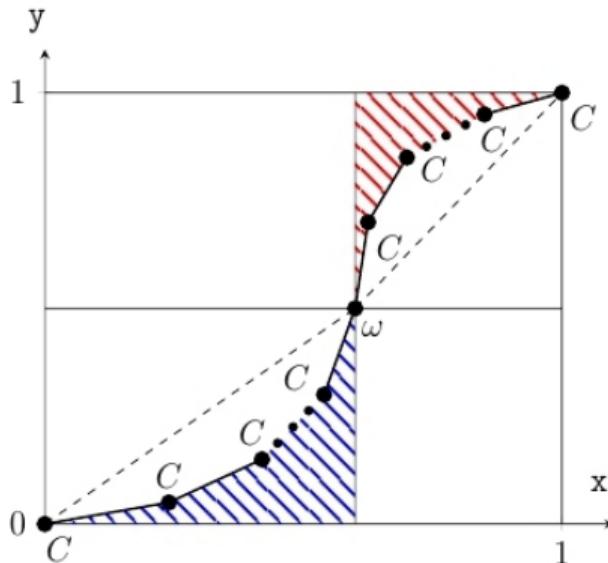
$$\mathcal{D}_n^{[n]} : \quad \eta_i \xi_{i+1} - \eta_{i+1} \xi_i \leq 0, \quad i = 1, \dots, n-1; \quad \frac{1}{2} S_n^{[n]} = \text{Area}$$



Geometric interpretation

ω in the k -th position:

$$\mathcal{D}_n^{[k]} = \begin{cases} \eta_i \xi_{i+1} - \eta_{i+1} \xi_i \leq 0, & i \in [1, k-1], \\ \eta_i \xi_{i+1} - \eta_{i+1} \xi_i \geq 0, & i \in [k+1, n-1]. \end{cases} \quad \frac{1}{2} S_n^k = \text{Area-Area}$$



Vertex dualities

Key property of the vertex Υ :

$$\Upsilon(\omega, C^n) \Big|_{y=0} = 0$$

Integrability constraint $d^2 = 0$ mixes vertices Υ and \mathcal{V}

$$\begin{aligned}\mathcal{V}_{\omega\omega C} * C - \omega * \Upsilon_{\omega CC} &\stackrel{y=0}{=} 0, & \mathcal{V}_{\omega C\omega} * C - \omega * \Upsilon_{C\omega C} &\stackrel{y=0}{=} 0, \\ \mathcal{V}_{C\omega\omega} * C - C * \pi(\mathcal{V}_{\omega\omega C}) &\stackrel{y=0}{=} 0, & -\omega * \Upsilon_{CC\omega} - \Upsilon_{\omega CC} * \pi(\omega) &\stackrel{y=0}{=} 0, \\ -\Upsilon_{C\omega C} * \pi(\omega) - C * \pi(\mathcal{V}_{\omega C\omega}) &\stackrel{y=0}{=} 0, & -\Upsilon_{CC\omega} * \pi(\omega) - C * \pi(\mathcal{V}_{C\omega\omega}) &\stackrel{y=0}{=} 0.\end{aligned}$$

In terms of sources:

$$\Psi_1^{[0,0]}(-p_2|t_1, t_2; p_1) = \Phi_2^{[0]}(t_1|t_2; p_1, p_2) \quad \rightarrow \quad \Psi_1^{[0,0]}(y|t_1, t_2; p) = \Phi_2^{[0]}(t_1|t_2; p, -y)$$

Vertex dualities

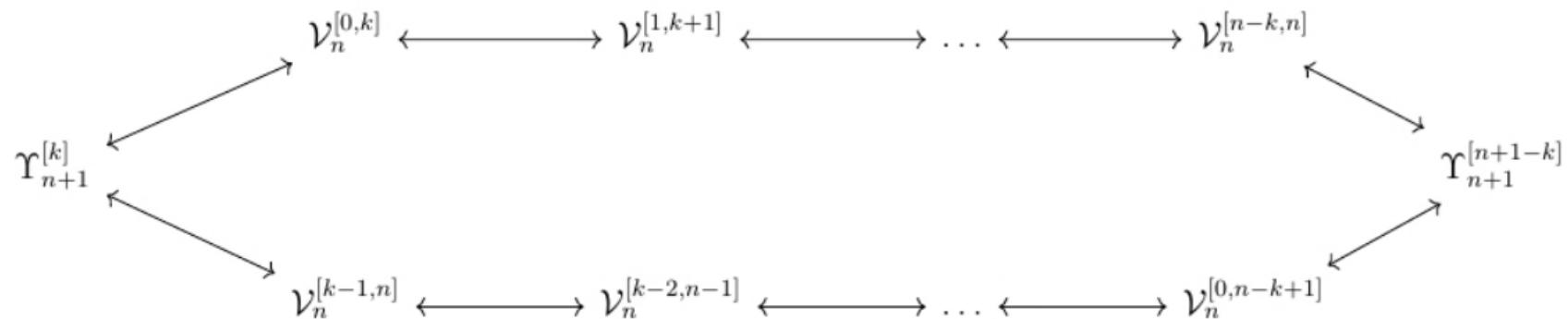
Lower-order example:

$$\begin{array}{ccc} \mathcal{V}_{\omega\omega C} & \longleftrightarrow & \mathcal{V}_{C\omega\omega} \\ \downarrow & & \downarrow \\ \Upsilon_{\omega CC} & \longleftrightarrow & \Upsilon_{CC\omega} \end{array}$$

$$\begin{array}{ccccc} & & \mathcal{V}_{\omega C\omega} & & \\ & \nearrow & & \swarrow & \\ \Upsilon_{C\omega C} & & & & \Upsilon_{C\omega C} \\ & \searrow & & \swarrow & \\ & & \mathcal{V}_{\omega C\omega} & & \end{array}$$

Vertex dualities

General case:



Interaction structure

Locality

Vertices Υ are spin-local:

$$\Upsilon(\omega, C^3) \sim \begin{array}{|c|c|c|c|} \hline & \bullet & \bullet & \bullet \\ \hline \bullet & & & \\ \hline \bullet & \bullet & & \\ \hline \end{array} \times \begin{array}{|c|c|c|c|} \hline & \bullet & \bullet & \bullet \\ \hline \bullet & & & \\ \hline \bullet & \bullet & & \\ \hline \end{array} \times \begin{array}{|c|c|c|c|} \hline & \bullet & \bullet & \bullet \\ \hline \bullet & & & \\ \hline \bullet & \bullet & & \\ \hline \end{array}$$

Finite number of contractions for fixed legs s_i

Vertices \mathcal{V} are spin ultra-local:

$$\mathcal{V}(\omega, \omega, C^n(0, \vec{y}))$$

Proliferated form and projectively-compact spin locality

$$d_x C_s^k = \sum_{i=0}^k \sum_{\vec{n}, \vec{s}} \Upsilon_s^i(\omega_{s'}^{n'}, \overbrace{C_{s_0}^{n_0} \dots C_{s_i}^{n_i}}^{i+1}), \quad C_s := C_s^0 \rightarrow C_s^1 \rightarrow \dots \rightarrow C_s^\infty = C_s^k, \quad k \geq 0$$

Conclusion

- All 'off-shell' bosonic HS vertices are manifestly found
- Vertices appear to be space-time spin local and minimal
- Condition $\Upsilon(\omega, C^n) \Big|_{y=0} = 0$ results in a net of vertex duality relations
- Precise form of the HS vertices reveals an interesting geometric structure: the integration phase space contains convex/concave polygons generalizing the Moyal star product geometric representation
- All-order exact unfolded equation for the Weyl module descendent C_s^k is a polynomial in C

The End